

Polynomial time approximation of entropy of shifts of finite type

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\mathbb{Z}^d Shifts of finite type

- Let \mathcal{A} be a finite alphabet.

$\mathcal{A}^{\mathbb{Z}^d} := \{ \text{all } d\text{-dimensional arrays of symbols from } \mathcal{A} \}.$

- **Shift of finite type (SFT):**

Let \mathcal{F} is a *finite* list of “forbidden” patterns on *finite* sets,

$$X = X_{\mathcal{F}} =$$

$\{x \in \mathcal{A}^{\mathbb{Z}^d} : x \text{ contains no translate of an element of } \mathcal{F}\}$

- SFT’s also known as “finite memory constraints.”
- **Nearest neighbor (n.n.) SFT:** an SFT where all forbidden patterns are patterns on *edges* of \mathbb{Z}^d .
- Main Example ($d = 2$): **hard square SFT**

$$\mathcal{A} = \{0, 1\}, \mathcal{F} = \left\{ \begin{array}{c} 11 \\ \downarrow \\ 1 \end{array} \right\}$$

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Topological entropy

- d -dimensional cube: $B_n := [0, n - 1]^d$
- for an SFT X ,

$$L_n(X) = \{ \text{legal configurations on } B_n \}$$

- **Topological entropy (noiseless capacity):**

$$h(X) := \lim_{n \rightarrow \infty} \frac{\log |L_n(X)|}{n^d}$$

- By subadditivity of $\log |L_n(X)|$,

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Golden Mean Shift ($(1, \infty)$ constraint): $\mathcal{F} = \{11\}$

- Adjacency matrix A of G is the square matrix indexed by \mathcal{A} :

$$A_{ab} = \begin{cases} 1 & ab \notin \mathcal{F} \\ 0 & ab \in \mathcal{F} \end{cases}$$

- $h(X) = \log \lambda(A)$, where $\lambda(A)$ is the spectral radius of A .
- Characterization of entropies for $d = 1$ (Lind):

$$\{\log \lambda^{1/q}\}$$

where λ is a Perron number and $q \in \mathbb{N}$

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Examples of \mathbb{Z}^2 SFTs: hard square

- **hard squares** $\mathcal{A} = \{0, 1\}, \mathcal{F} = \left\{ \begin{array}{c} 11 \\ 1 \end{array} \right\}$
- $h(\text{hard square SFT}) = ???$
- (Baxter) $h(\text{hard hexagons}) = \log(\lambda)$ where λ is an algebraic integer of degree 24.

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Examples of \mathbb{Z}^2 SFTs: checkerboard (coloring) constraints

- **q -checkerboard \mathcal{C}_q :** $\mathcal{A} = \{1, \dots, q\}, \mathcal{F} = \{aa, \begin{smallmatrix} a \\ a \end{smallmatrix}\}$
- $h(\mathcal{C}_2) = 0$
- (Lieb): $h(\mathcal{C}_3) = (3/2) \log(4/3)$
- $h(\mathcal{C}_4) = ???$

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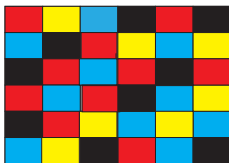
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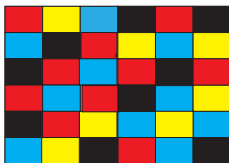
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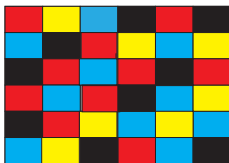
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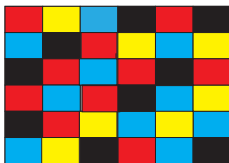
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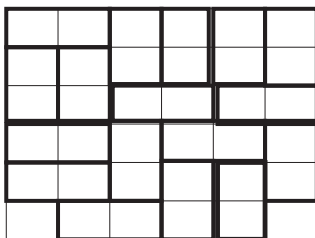
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Examples of \mathbb{Z}^2 SFT's: dimers

- **dimers:**



$$\mathcal{F} = \{LL, LT, LB, RR, TR, BR, \begin{matrix} T & T & T & B & L & R \\ L & R & T & B & B & B \end{matrix}\}$$

- (Fisher-Kastelyn-Temperley):

$$h(\text{Dimers}) = \frac{1}{16\pi^2} \int_{-\pi}^{\pi} \int_{-\pi}^{\pi} \log(4 + 2\cos\theta + 2\cos\phi) d\theta d\phi$$

- $h(\text{Monomers-Dimers}) = ???$

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L	R	T	T	T	T
T	T	B	B	B	B
B	B	L	R	L	R
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Topological entropy, $d \geq 2$

- Exact formula known only in a few cases.
- Characterization of entropies for $d \geq 2$
(Hochman-Meyerovitch):

{**right recursively enumerable** (RRE) numbers $h \geq 0$ }

i.e, there is an algorithm that produces a sequence $r_n \geq h$
s.t. $r_n \rightarrow h$.

Proof:

- Necessity: Let $r_n := \frac{\log |L_n|}{n^d}$.
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Since $\lim = \inf$, each $r_n \geq h$.
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- Exact formula known only in a few cases.
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Given a shift-invariant Borel probability measure μ on $\mathcal{A}^{\mathbb{Z}^d}$,

- For finite $S \in \mathbb{Z}^d$,

$$H_\mu(S) := \sum_{x \in \mathcal{A}^S} -\mu(x) \log \mu(x) = \int -\log \mu(x) d\mu(x)$$

- For finite disjoint S, T ,

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- $h(\mu) := \lim_{n \rightarrow \infty} \frac{H_\mu(B_n)}{n^d}$
- $d = 1$: Theorem: $h(\mu) = H_\mu(0 \mid \{-1, -2, -3, \dots\})$
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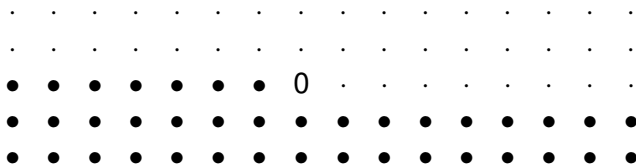


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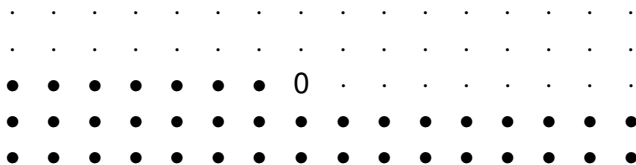


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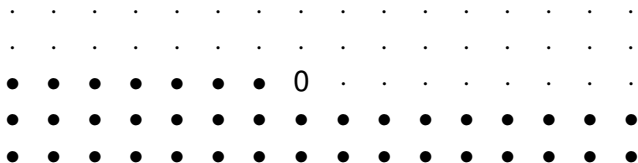
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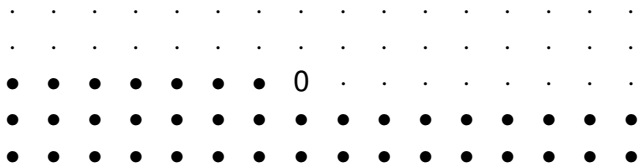
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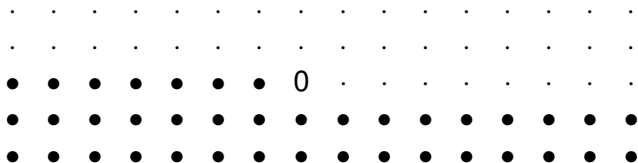
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Variational Principle for Topological Entropy

- For an SFT X ,

$$h(X) = \sup_{\mu} h(\mu)$$

where the sup is taken over all shift-invariant Borel probability measures μ s.t. $\text{support}(\mu) \subseteq X$.

- Fact: The sup is always achieved. A measure which achieves the sup is called a **measure of maximal entropy (MME)**.
- So for an MME μ , $h(X) = h(\mu) = \int I_{\mu}(x) d\mu(x)$
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This is an average of n^d terms of two types:

- *Bulk terms*: Terms that are far from the boundary of B_n
- *Boundary terms*: Terms that are near the boundary of B_n

Bulk terms are close to $I_\mu(s^{\mathbb{Z}^d})$. All terms are uniformly bounded. Most terms are bulk terms. So, $h(X) = I_\mu(s^{\mathbb{Z}^d})$

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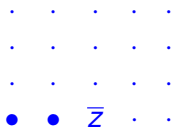
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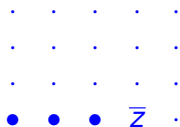
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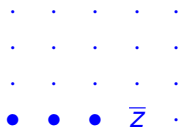
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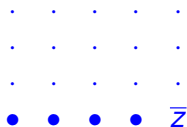
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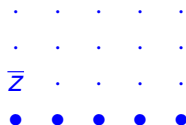
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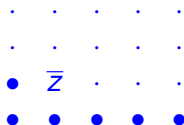
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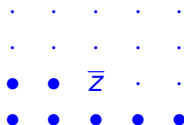
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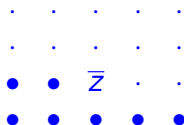
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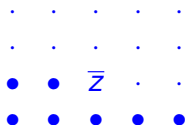
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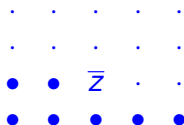
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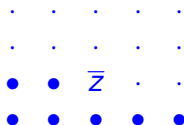
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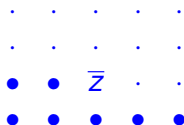
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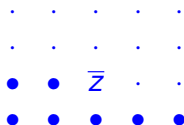
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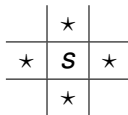


Examples: Yes: Hard squares ($s = 0$)

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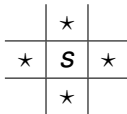


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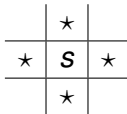


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Example: $R_{3,4,3}$:



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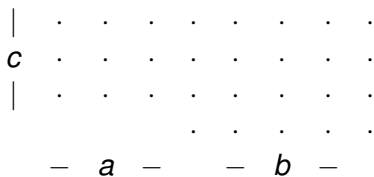
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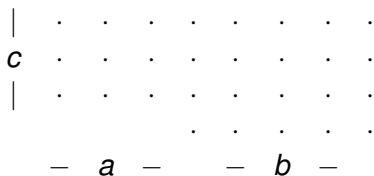
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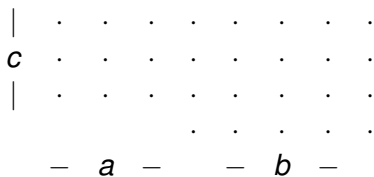
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Moreover, if $d = 2$ and convergence in hypothesis 2 is exponential, then there is a polynomial time algorithm to compute $h(X)$.

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- Accuracy is $e^{-\Omega(n)}$
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Proof of Claim, via transfer matrices

$$\mu(s^0 \mid s^{\partial R_{n,n,n}}) = \frac{\begin{array}{cccccc} & s & s & s & s & s \\ s & \cdot & \cdot & \cdot & \cdot & \cdot & s \\ \# & s & \cdot & \cdot & \cdot & \cdot & s \\ s & s & s & s & \cdot & \cdot & s \\ & & & s & s & s \\ s & s & s & s & s & s \\ s & \cdot & \cdot & \cdot & \cdot & \cdot & s \\ \# & s & \cdot & \cdot & \cdot & \cdot & s \\ s & s & s & \cdot & \cdot & \cdot & s \\ & & & s & s & s \end{array}}{\begin{array}{cccccc} & s & s & s & s & s \\ s & \cdot & \cdot & \cdot & \cdot & \cdot & s \\ \# & s & \cdot & \cdot & \cdot & \cdot & s \\ s & s & s & \cdot & \cdot & \cdot & s \\ & & & s & s & s \end{array}}$$

$$= \frac{(\prod_{i=-n}^{-1} M_i) \hat{M}_0 (\prod_{i=1}^{n-1} M_i)}{(\prod_{i=-n}^{-1} M_i) M_0 (\prod_{i=1}^{n-1} M_i)}$$

M_i is transition matrix from column i to column $i + 1$ compatible with $s^{\partial R_{n,n,n}}$ and

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- Weaken fixed point $s^{\mathbb{Z}^d}$ to periodic orbit
- Weaken safe symbol to topological strong spatial mixing
- Applies to
 - hard squares
 - monomer-dimers
 - q -checkerboard SFT with $q \geq 6$
- Generalize results from entropy to *pressure* of n n. interactions on n .n. SFT's
- Applies to large sets of temperature regions for classical models in statistical physics, in both subcritical and supercritical regions:
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End of talk

The following slides form a hodge-podge of topics that were not included in the talk.

Topological Strong Spatial Mixing (TSSM)

Defn of TSSM with gap g :

for any disjoint $U, S, V \in Z^d$ s.t. $d(U, V) \geq g$,

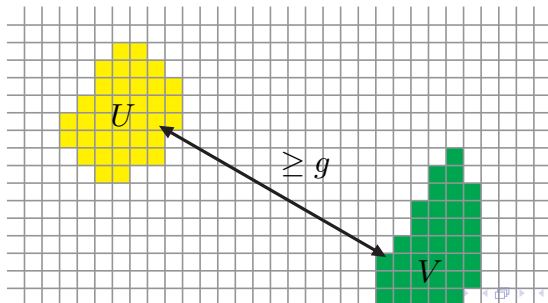
if $u \in A^U$, $s \in A^S$, $v \in A^V$, s.t. us and sv are allowed, then so is usv .

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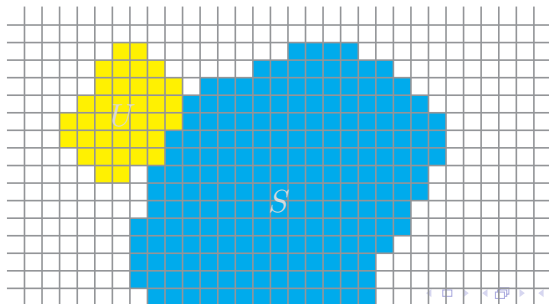


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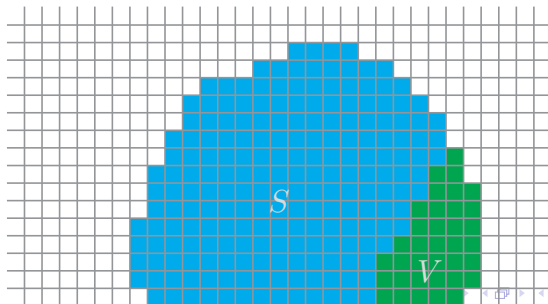


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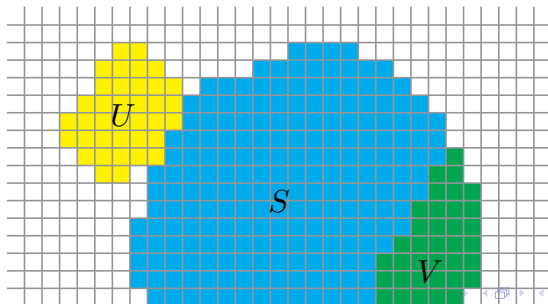
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Let $R_{a,b,c} := [-a, -1] \times [1, c] \cup [0, b] \times [0, c]$

Example: $R_{3,4,3}$:



Theorem: Let X be a \mathbb{Z}^d n.n. SFT and μ an MME on X . If

- 1 X satisfies TSSM
- 2 (for $d = 2$) For some periodic orbit O in X and all $\omega \in O$

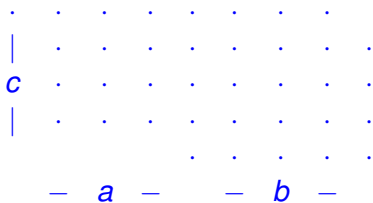
$$L(\omega) := \lim_{a,b,c \rightarrow \infty} \mu(\omega(0) \mid \omega(\partial R_{a,b,c})) \text{ exists}$$

Then

$$h(X) = -\frac{1}{|O|} \sum_{\omega \in O} \log L(\omega)$$

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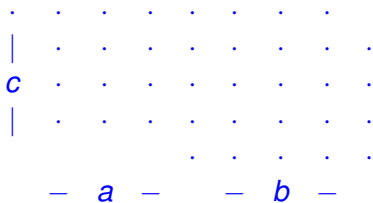
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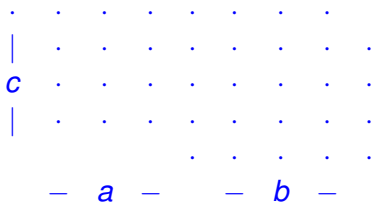
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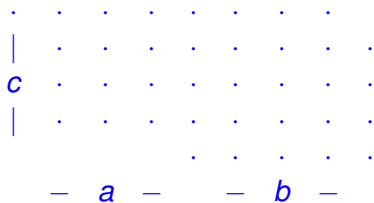
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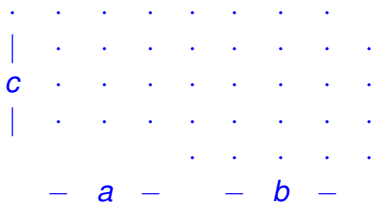
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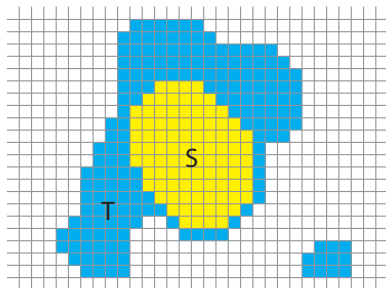
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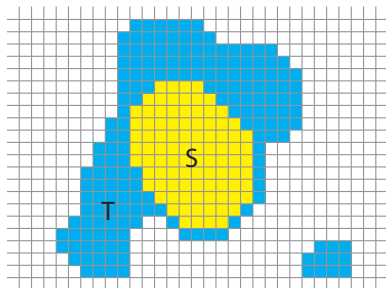
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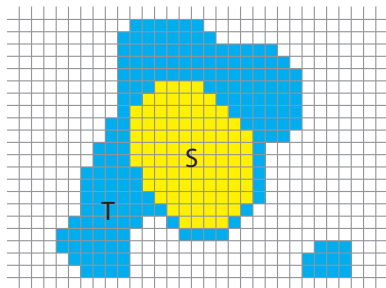
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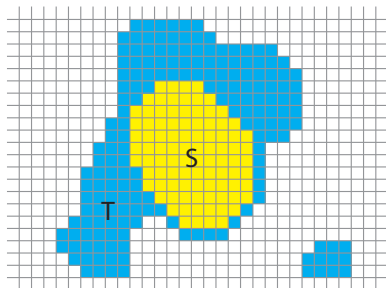
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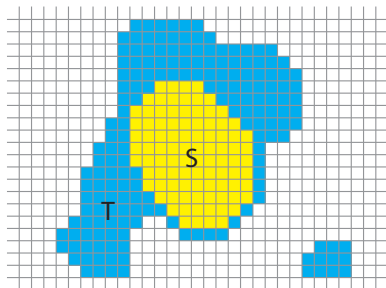
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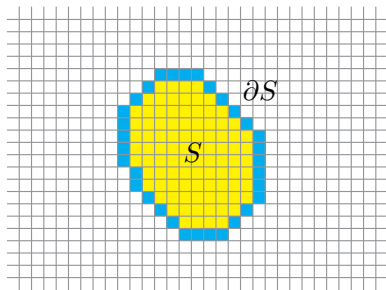
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$$L_S^y(X) := \{x \in \mathcal{A}^S : xy \text{ is legal} \}$$

An MRF on X is **uniform** if whenever $\mu(y) > 0$, then for $x \in L_S^y(X)$

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Theorem (Lanford/Ruelle, Burton/Steif): Every MME on a n.n. SFT is a uniform MRF.

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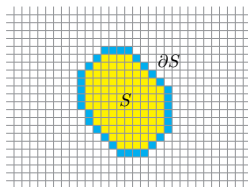
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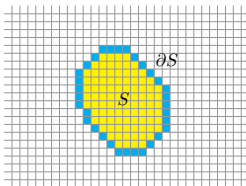
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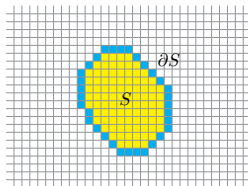
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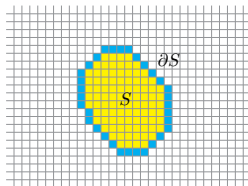
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$$\mu(x \mid y) = \frac{1}{|L_S^y(X)|}$$

Theorem (Lanford/Ruelle, Burton/Steif): Every MME on a n.n. SFT is a uniform MRF.

- Since μ is an MME, μ must be a uniform MRF.
- Since s is a safe symbol,
 - 1 For all $T \in \mathbb{Z}^d$ containing 0,

$$\mu(s^0 \mid s^{\partial T}) \geq \frac{1}{|\mathcal{A}|}.$$

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$$h(X) = \lim_{n \rightarrow \infty} \frac{-\log \mu(s^{B_n} \mid s^{\partial B_n})}{n^d}$$

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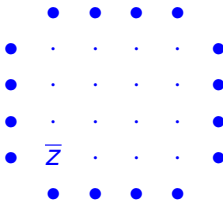
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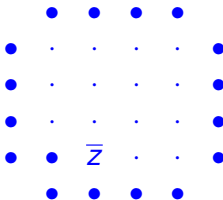
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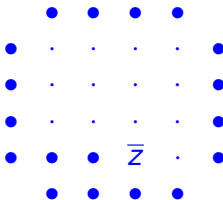
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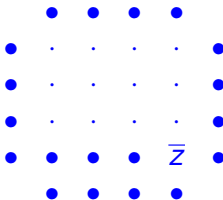
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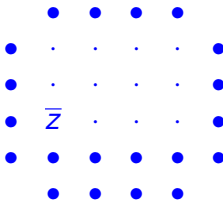
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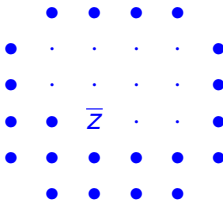
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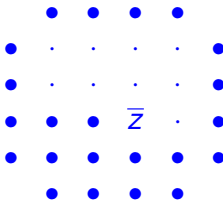
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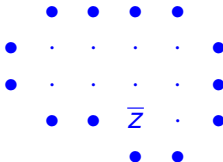
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Algorithmic consequence

Theorem: Let X be a n.n. \mathbb{Z}^2 SFT and μ an MME on X . If

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Strong Spatial Mixing

- An MRF μ satisfies **strong spatial mixing (SSM)** at rate $f(n)$

if for all $V \in Z^d$, $U \subset V$

all $u \in A^U$, and $v, v' \in A^{\partial V}$ satisfying $\mu(v), \mu(v') > 0$,

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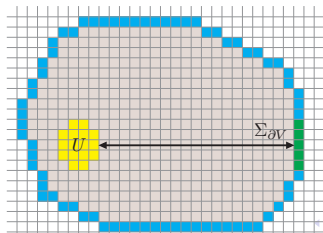
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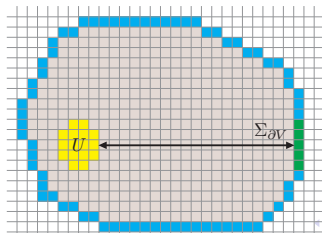
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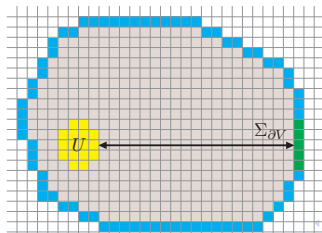
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Stronger conclusion

Theorem (Briceno): Let X be a \mathbb{Z}^d n.n. SFT and μ an MME on X . If

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Then for *all* invariant measures ν s.t. $\text{support}(\nu) \subseteq X$,

$$h(X) = \int I_{\mu}(x) d\nu(x)$$

Applies to:

- hard squares
- q -checkerboard with $q \geq 6$

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Topological Pressure and Variational Principle

- Let X be a shift space and $f : X \rightarrow \mathbb{R}$ a continuous function.
- **Topological Pressure** (defined by Variational Principle):

$$P_X(f) := \sup_{\mu} h(\mu) + \int f d\mu$$

where the sup is taken over all shift-invariant Borel probability measures μ such that $\text{support}(\mu) \subseteq X$.

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Nearest-Neighbour interactions and Gibbs measures

- A *nearest-neighbor interaction* is a shift-invariant function Φ from a set of configurations on vertices and edges in \mathbb{Z}^d to $\mathbb{R} \cup \infty$
- For a nearest-neighbor interaction Φ , the *underlying SFT*:
$$X = X_\Phi := \{x \in \mathcal{A}^{\mathbb{Z}^d} : \Phi(x(\{v, v'\})) \neq \infty, \text{ for all } v \sim v'\}.$$
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Examples of n.n. Gibbs measures

- uniform MME on n.n. SFT
- hard square model with activities
- ferromagnetic Ising model with no external field.

Equilibrium states versus n.n. Gibbs measures

- Pressure of n.n. interaction Φ :

$$P(\Phi) := \lim_{n \rightarrow \infty} \frac{\log Z^\Phi(B_n)}{n^d}$$

where $Z^\Phi(B_n)$ is the “free boundary” normalization.

- Let $A_\Phi(x) := -\Phi(x(0)) - \sum_{i=1}^d \Phi(x(0), x(e_i))$.
- Fact: $P_{X_\Phi}(A_\Phi) = P(\Phi)$.
- Lanford-Ruelle Theorem: Every equilibrium state for A_Φ is a Gibbs measure for Φ .
- Dobrushin Theorem: If X_Φ is strongly irreducible, then every Gibbs measure for Φ is an equilibrium state for A_Φ .
- These theorems hold in much greater generality.

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Pressure representation and approximation

Theorem (Adams, Briceno, Marcus, Pavlov): Let μ a Gibbs measure for a n.n. interaction Φ with underlying \mathbb{Z}^d n.n. SFT X .
If

- 1 X satisfies TSSM
- 2 For some periodic orbit O in X and all $\omega \in O$

$$L(\omega) := \lim_{a,b,c \rightarrow \infty} \mu(\omega(0) \mid \omega(\partial R_{a,b,c})) \text{ exists}$$

Then

$$P(\Phi) = \frac{1}{|O|} \sum_{\omega \in O} -\log L(\omega) + A_{\Phi}(\omega)$$

Moreover, if $d = 2$ and convergence in hypothesis 2 is exponential, then there is a polynomial time algorithm to compute $P(\Phi)$.

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Stronger conclusion

Theorem (Briceno): Let μ a Gibbs measure for a n.n. interaction Φ with underlying \mathbb{Z}^d n.n. SFT X . If

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Then for all shift-invariant measures ν such that $\text{support}(\nu) \subseteq X$,

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An SFT X satisfies the **D-condition** if

- there exist sequences of finite subsets $(\Lambda_n), (M_n)$ of \mathbb{Z}^d such that $\Lambda_n \nearrow \infty$, $\Lambda_n \subseteq M_n$, $\frac{|M_n|}{|\Lambda_n|} \rightarrow 1$, such that
- for any globally admissible $v \in \mathcal{A}^{\Lambda_n}$ and finite $S \subset M_n^c$ and globally admissible $w \in \mathcal{A}^S$, we have that vw is globally admissible.

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Connection with Thermodynamic Formalism

Theorem: Let μ a Gibbs measure for a n.n. interaction Φ with underlying \mathbb{Z}^d n.n. SFT X . If

- X satisfies the D-condition
- $I_\mu = A_\Psi$ for some *absolutely summable* interaction Ψ s.t.
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MME, $d = 1$

- Assuming adjacency matrix A is irreducible and aperiodic, there is a unique MME μ_{\max} , which is a Markov chain given by transition matrix

$$P_{ij} = \begin{cases} \frac{r_j}{\lambda r_i} & ij \notin \mathcal{F} \\ 0 & ij \in \mathcal{F} \end{cases}$$

where $\lambda = \lambda(A)$ and r is a right eigenvector for λ , and stationary vector $r_i \ell_j$ where ℓ is a left eigenvector for λ (suitably normalized)

- Thus, if $\mu(w_1 w_2 \dots w_{n-1} w_n) > 0$, then

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Entropy representation for MME, $d = 1$



$$\begin{aligned}I_{\mu}(x) &= -\log \mu(x(0) | x(\mathcal{P})) \\ &= -\log P_{x_0 x_{-1}} \\ &= \log \lambda + \log r_{x_{-1}} - \log r_{x_0}\end{aligned}$$

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In particular, if the SFT has a fixed point $x^* := a^{\mathbb{Z}}$ and ν is the delta measure on x^* , then on

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