

Yamabe flow and its soliton on manifolds with boundary

Pak Tung Ho, Tamkang University

Analysis and PDE Seminar at HKU

21st of November, 2023

Joint work with Jinwoo Shin at Korea Institute for Advanced Study
(KIAS)

Yamabe problem

Suppose (M, g_0) is an n -dimensional closed (i.e. compact without boundary) Riemannian manifold, where $n \geq 3$.

Yamabe problem

Suppose (M, g_0) is an n -dimensional closed (i.e. compact without boundary) Riemannian manifold, where $n \geq 3$.

The **Yamabe problem** is to find a metric g conformal to g_0 such that its scalar curvature R_g is constant.

Yamabe problem

Suppose (M, g_0) is an n -dimensional closed (i.e. compact without boundary) Riemannian manifold, where $n \geq 3$.

The **Yamabe problem** is to find a metric g conformal to g_0 such that its scalar curvature R_g is constant.

If we write $g = u^{\frac{4}{n-2}} g_0$, then

$$-\frac{4(n-1)}{n-2} \Delta_{g_0} u + R_{g_0} u = R_g u^{\frac{n+2}{n-2}}. \quad (\text{YamabePDE})$$

Yamabe problem

Suppose (M, g_0) is an n -dimensional closed (i.e. compact without boundary) Riemannian manifold, where $n \geq 3$.

The **Yamabe problem** is to find a metric g conformal to g_0 such that its scalar curvature R_g is constant.

If we write $g = u^{\frac{4}{n-2}} g_0$, then

$$-\frac{4(n-1)}{n-2} \Delta_{g_0} u + R_{g_0} u = R_g u^{\frac{n+2}{n-2}}. \quad (\text{YamabePDE})$$

The Yamabe problem is to find $0 < u \in C^\infty(M)$ with $R_g \equiv c$.

Yamabe problem

The **Yamabe constant** of (M, g_0) is

$$Y(M, g_0) = \inf \{ E_{g_0}(u) : 0 < u \in C^\infty(M) \}.$$

Here

$$E_{g_0}(u) = \frac{\int_M \frac{4(n-1)}{n-2} |\nabla_{g_0} u|^2 + R_{g_0} u^2 dV_{g_0}}{\left(\int_M u^{\frac{2n}{n-2}} dV_{g_0} \right)^{\frac{n-2}{n}}} = \frac{\int_M R_g dV_g}{\left(\int_M dV_g \right)^{\frac{n-2}{n}}}$$

is the Yamabe energy, where $g = u^{\frac{4}{n-2}} g_0$.

Yamabe problem

The **Yamabe constant** of (M, g_0) is

$$Y(M, g_0) = \inf \{ E_{g_0}(u) : 0 < u \in C^\infty(M) \}.$$

Here

$$E_{g_0}(u) = \frac{\int_M \frac{4(n-1)}{n-2} |\nabla_{g_0} u|^2 + R_{g_0} u^2 dV_{g_0}}{\left(\int_M u^{\frac{2n}{n-2}} dV_{g_0} \right)^{\frac{n-2}{n}}} = \frac{\int_M R_g dV_g}{\left(\int_M dV_g \right)^{\frac{n-2}{n}}}$$

is the Yamabe energy, where $g = u^{\frac{4}{n-2}} g_0$.

Fact: If u is a positive minimizer, i.e. $E_{g_0}(u) = Y(M, g_0)$, then u solves (YamabePDE).

Proof: $\frac{d}{dt} E_{g_0}(u + tv) = 0$ for all $v \in C^\infty(M)$.

Yamabe problem

The **Yamabe constant** of (M, g_0) is

$$Y(M, g_0) = \inf \{ E_{g_0}(u) : 0 < u \in C^\infty(M) \}.$$

Here

$$E_{g_0}(u) = \frac{\int_M \frac{4(n-1)}{n-2} |\nabla_{g_0} u|^2 + R_{g_0} u^2 dV_{g_0}}{\left(\int_M u^{\frac{2n}{n-2}} dV_{g_0} \right)^{\frac{n-2}{n}}} = \frac{\int_M R_g dV_g}{\left(\int_M dV_g \right)^{\frac{n-2}{n}}}$$

is the Yamabe energy, where $g = u^{\frac{4}{n-2}} g_0$.

Fact: If u is a positive minimizer, i.e. $E_{g_0}(u) = Y(M, g_0)$, then u solves (YamabePDE).

Proof: $\frac{d}{dt} E_{g_0}(u + tv) = 0$ for all $v \in C^\infty(M)$.

The Yamabe problem was solved by Aubin, Trudinger, and Schoen.

Hamilton introduced the **Yamabe flow**:

$$\frac{\partial}{\partial t} g(t) = -(R_{g(t)} - \bar{R}_{g(t)})g(t), \quad g(0) = g_0, \quad (\text{YF})$$

where $\bar{R}_{g(t)}$ is the average of $R_{g(t)}$:

$$\bar{R}_{g(t)} = \frac{\int_M R_{g(t)} dV_{g(t)}}{\int_M dV_{g(t)}}.$$

Hamilton introduced the **Yamabe flow**:

$$\frac{\partial}{\partial t} g(t) = -(R_{g(t)} - \bar{R}_{g(t)})g(t), \quad g(0) = g_0, \quad (\text{YF})$$

where $\bar{R}_{g(t)}$ is the average of $R_{g(t)}$:

$$\bar{R}_{g(t)} = \frac{\int_M R_{g(t)} dV_{g(t)}}{\int_M dV_{g(t)}}.$$

Hamilton proved the long time existence of (YF). Also, he proved the convergence of (YF) when $Y(M, g_0) \leq 0$, i.e. $g(t) \rightarrow g_\infty$ as $t \rightarrow \infty$ for some metric g_∞ with $R_{g_\infty} \equiv c$.

Hamilton introduced the **Yamabe flow**:

$$\frac{\partial}{\partial t} g(t) = -(R_{g(t)} - \bar{R}_{g(t)})g(t), \quad g(0) = g_0, \quad (\text{YF})$$

where $\bar{R}_{g(t)}$ is the average of $R_{g(t)}$:

$$\bar{R}_{g(t)} = \frac{\int_M R_{g(t)} dV_{g(t)}}{\int_M dV_{g(t)}}.$$

Hamilton proved the long time existence of (YF). Also, he proved the convergence of (YF) when $Y(M, g_0) \leq 0$, i.e. $g(t) \rightarrow g_\infty$ as $t \rightarrow \infty$ for some metric g_∞ with $R_{g_\infty} \equiv c$.

When $Y(M, g_0) > 0$, the convergence of (YF) was studied by Brendle, Chow, Schwetlick and Struwe, and Ye.

Theorem (Carlotto-Chodosh-Rubinstein)

If g_∞ is an integrable critical point of Yamabe energy, then (YF) converges exponentially to g_∞ , i.e.

$$\|g(t) - g_\infty\|_{C^{2,\alpha}(M, g_\infty)} \leq Ce^{-\delta t}$$

for some $\delta > 0$.

Theorem (Carlotto-Chodosh-Rubinstein)

If g_∞ is an integrable critical point of Yamabe energy, then (YF) converges exponentially to g_∞ , i.e.

$$\|g(t) - g_\infty\|_{C^{2,\alpha}(M, g_\infty)} \leq Ce^{-\delta t}$$

for some $\delta > 0$.

Fact: g_∞ is an integrable critical point of Yamabe energy if $\ker \mathcal{L}_\infty = \{0\}$, where

$$\mathcal{L}_\infty v = (n-1)\Delta_{g_\infty} v + R_{g_\infty} v$$

is the linearized Yamabe operator of g_∞ .

Theorem (Carlotto-Chodosh-Rubinstein)

There exists g_∞ such that (YF) does not converge exponentially to g_∞ . More precisely,

$$C^{-1}(1+t)^{-\frac{1}{p-2}} \leq \|g(t) - g_\infty\|_{C^{2,\alpha}(M, g_\infty)} \leq C(1+t)^{-\frac{1}{p-2}}$$

for $p = 3$ and for some $p \geq 3$.

Theorem (Carlotto-Chodosh-Rubinstein)

There exists g_∞ such that (YF) does not converge exponentially to g_∞ . More precisely,

$$C^{-1}(1+t)^{-\frac{1}{p-2}} \leq \|g(t) - g_\infty\|_{C^{2,\alpha}(M, g_\infty)} \leq C(1+t)^{-\frac{1}{p-2}}$$

for $p = 3$ and for some $p \geq 3$.

More precisely, Carlotto-Chodosh-Rubinstein found metrics g_∞ which satisfy the AS_p .

Theorem (Carlotto-Chodosh-Rubinstein)

There exists g_∞ such that (YF) does not converge exponentially to g_∞ . More precisely,

$$C^{-1}(1+t)^{-\frac{1}{p-2}} \leq \|g(t) - g_\infty\|_{C^{2,\alpha}(M, g_\infty)} \leq C(1+t)^{-\frac{1}{p-2}}$$

for $p = 3$ and for some $p \geq 3$.

More precisely, Carlotto-Chodosh-Rubinstein found metrics g_∞ which satisfy the AS_p .

Fact: g_∞ satisfies AS_3 if there exists $v \neq 0$ such that

$$\mathcal{L}_\infty v = (n-1)\Delta_{g_\infty} v + R_{g_\infty} v = 0 \quad \text{and} \quad R_{g_\infty} \int_M v^3 dV_{g_\infty} \neq 0.$$

Theorem (Carlotto-Chodosh-Rubinstein)

There exists g_∞ such that (YF) does not converge exponentially to g_∞ . More precisely,

$$C^{-1}(1+t)^{-\frac{1}{p-2}} \leq \|g(t) - g_\infty\|_{C^{2,\alpha}(M, g_\infty)} \leq C(1+t)^{-\frac{1}{p-2}}$$

for $p = 3$ and for some $p \geq 3$.

More precisely, Carlotto-Chodosh-Rubinstein found metrics g_∞ which satisfy the AS_p .

Fact: g_∞ satisfies AS_3 if there exists $v \neq 0$ such that

$$\mathcal{L}_\infty v = (n-1)\Delta_{g_\infty} v + R_{g_\infty} v = 0 \quad \text{and} \quad R_{g_\infty} \int_M v^3 dV_{g_\infty} \neq 0.$$

Goal: Find g_∞ with $R_{g_\infty} \equiv c > 0$ such that $\int_M v^3 dV_{g_\infty} \neq 0$.

Yamabe problem with boundary

Suppose (M, g_0) is an n -dimensional compact Riemannian manifold with boundary ∂M , where $n \geq 3$.

Yamabe problem with boundary

Suppose (M, g_0) is an n -dimensional compact Riemannian manifold with boundary ∂M , where $n \geq 3$.

There are two types of **Yamabe problem with boundary**:

(I) Find g conformal to g_0 such that

$$R_g = c \text{ in } M \quad \text{and} \quad H_g = 0 \text{ on } \partial M.$$

Yamabe problem with boundary

Suppose (M, g_0) is an n -dimensional compact Riemannian manifold with boundary ∂M , where $n \geq 3$.

There are two types of **Yamabe problem with boundary**:

(I) Find g conformal to g_0 such that

$$R_g = c \text{ in } M \quad \text{and} \quad H_g = 0 \text{ on } \partial M.$$

(II) Find g conformal to g_0 such that

$$R_g = 0 \text{ in } M \quad \text{and} \quad H_g = c \text{ on } \partial M.$$

Yamabe problem with boundary

Suppose (M, g_0) is an n -dimensional compact Riemannian manifold with boundary ∂M , where $n \geq 3$.

There are two types of **Yamabe problem with boundary**:

(I) Find g conformal to g_0 such that

$$R_g = c \text{ in } M \quad \text{and} \quad H_g = 0 \text{ on } \partial M.$$

(II) Find g conformal to g_0 such that

$$R_g = 0 \text{ in } M \quad \text{and} \quad H_g = c \text{ on } \partial M.$$

$(\mathbb{S}_+^n, g_{\mathbb{S}_+^n})$ is an example of (i), and (D^n, g_{flat}) is an example of (ii).

Yamabe problem with boundary

If $g = u^{\frac{4}{n-2}} g_0$, then

$$-\frac{4(n-1)}{n-2} \Delta_{g_0} u + R_{g_0} u = R_g u^{\frac{n+2}{n-2}} \text{ in } M,$$
$$\frac{2(n-1)}{n-2} \frac{\partial u}{\partial \nu_{g_0}} + H_{g_0} u = H_g u^{\frac{n}{n-2}} \text{ on } \partial M.$$

Yamabe problem with boundary

If $g = u^{\frac{4}{n-2}} g_0$, then

$$-\frac{4(n-1)}{n-2} \Delta_{g_0} u + R_{g_0} u = R_g u^{\frac{n+2}{n-2}} \text{ in } M,$$
$$\frac{2(n-1)}{n-2} \frac{\partial u}{\partial \nu_{g_0}} + H_{g_0} u = H_g u^{\frac{n}{n-2}} \text{ on } \partial M.$$

The Yamabe problem with boundary was first introduced and studied by Escobar.

Yamabe problem with boundary

If $g = u^{\frac{4}{n-2}} g_0$, then

$$-\frac{4(n-1)}{n-2} \Delta_{g_0} u + R_{g_0} u = R_g u^{\frac{n+2}{n-2}} \text{ in } M,$$
$$\frac{2(n-1)}{n-2} \frac{\partial u}{\partial \nu_{g_0}} + H_{g_0} u = H_g u^{\frac{n}{n-2}} \text{ on } \partial M.$$

The Yamabe problem with boundary was first introduced and studied by Escobar.

The Yamabe problem with boundary was later studied by Almaraz, S. Chen, Marques, Mayer and Ndiaye, etc.

Yamabe flow with boundary

One wants to study the Yamabe problem with boundary by using geometric flow.

Yamabe flow with boundary

One wants to study the Yamabe problem with boundary by using geometric flow.

To study (I), we consider the **Yamabe flow with boundary**:

$$\begin{cases} \frac{\partial}{\partial t} g(t) = -(R_{g(t)} - \bar{R}_{g(t)})g(t) & \text{in } M, \\ H_{g(t)} = 0 \text{ on } \partial M, \quad g(0) = g_0. \end{cases} \quad (\text{YFB})$$

Yamabe flow with boundary

One wants to study the Yamabe problem with boundary by using geometric flow.

To study (I), we consider the **Yamabe flow with boundary**:

$$\begin{cases} \frac{\partial}{\partial t} g(t) = -(R_{g(t)} - \bar{R}_{g(t)})g(t) & \text{in } M, \\ H_{g(t)} = 0 \text{ on } \partial M, \quad g(0) = g_0. \end{cases} \quad (\text{YFB})$$

To study (II), we consider the **conformal mean curvature flow**:

$$\begin{cases} \frac{\partial}{\partial t} g(t) = -(H_{g(t)} - \bar{H}_{g(t)})g(t) & \text{on } \partial M, \\ R_{g(t)} = 0 \text{ in } M, \quad g(0) = g_0. \end{cases} \quad (\text{CMCF})$$

Yamabe flow with boundary

- ▶ (YFB) and (CMCF) are first introduced by Brendle, who proved the short-time existence. He also proved the convergence when the Yamabe constant is nonpositive.

Yamabe flow with boundary

- ▶ (YFB) and (CMCF) are first introduced by Brendle, who proved the short-time existence. He also proved the convergence when the Yamabe constant is nonpositive.
- ▶ When the Yamabe constant is positive, the convergence of (YFB) is proved by Almaraz-Sun, and the convergence of (CMCF) is proved by Almaraz.

Theorem (H.-Shin)

If g_∞ is an integrable critical point of Yamabe energy, then (YFB) (respectively (CMCF)) converges exponentially to g_∞ , i.e.

$$\|g(t) - g_\infty\|_{C^{2,\alpha}(M, g_\infty)} \leq Ce^{-\delta t}$$

for some $\delta > 0$.

Theorem (H.-Shin)

If g_∞ is an integrable critical point of Yamabe energy, then (YFB) (respectively (CMCF)) converges exponentially to g_∞ , i.e.

$$\|g(t) - g_\infty\|_{C^{2,\alpha}(M, g_\infty)} \leq Ce^{-\delta t}$$

for some $\delta > 0$.

Theorem (H.-Shin)

There exists g_∞ such that (YFB) (respectively (CMCF)) does not converge exponentially to g_∞ . More precisely,

$$C^{-1}(1+t)^{-\frac{1}{p-2}} \leq \|g(t) - g_\infty\|_{C^{2,\alpha}(M, g_\infty)} \leq C(1+t)^{-\frac{1}{p-2}}$$

for $p = 3$.

Sketlov eigenvalue

For a Riemannian manifold (M, g) with boundary ∂M , the **Dirichlet-to-Neumann map** $DN : C^\infty(\partial M) \rightarrow C^\infty(\partial M)$ is given by

$$DN(f) = \frac{\partial \hat{f}}{\partial \nu_g},$$

where \hat{f} is the harmonic extension of f ,

Sketlov eigenvalue

For a Riemannian manifold (M, g) with boundary ∂M , the **Dirichlet-to-Neumann map** $DN : C^\infty(\partial M) \rightarrow C^\infty(\partial M)$ is given by

$$DN(f) = \frac{\partial \hat{f}}{\partial \nu_g},$$

where \hat{f} is the harmonic extension of f , i.e.

$$\Delta_g \hat{f} = 0 \text{ in } M \quad \text{and} \quad \hat{f} = f \text{ on } \partial M.$$

Steklov eigenvalue

For a Riemannian manifold (M, g) with boundary ∂M , the **Dirichlet-to-Neumann map** $DN : C^\infty(\partial M) \rightarrow C^\infty(\partial M)$ is given by

$$DN(f) = \frac{\partial \hat{f}}{\partial \nu_g},$$

where \hat{f} is the harmonic extension of f , i.e.

$$\Delta_g \hat{f} = 0 \text{ in } M \quad \text{and} \quad \hat{f} = f \text{ on } \partial M.$$

The eigenvalue of DN is called **Steklov eigenvalue** of M ,

Steklov eigenvalue

For a Riemannian manifold (M, g) with boundary ∂M , the **Dirichlet-to-Neumann map** $DN : C^\infty(\partial M) \rightarrow C^\infty(\partial M)$ is given by

$$DN(f) = \frac{\partial \hat{f}}{\partial \nu_g},$$

where \hat{f} is the harmonic extension of f , i.e.

$$\Delta_g \hat{f} = 0 \text{ in } M \quad \text{and} \quad \hat{f} = f \text{ on } \partial M.$$

The eigenvalue of DN is called **Steklov eigenvalue** of M , i.e. σ is Steklov eigenvalue if

$$\Delta_g f = 0 \text{ in } M \quad \text{and} \quad \frac{\partial f}{\partial \nu_g} = \sigma f \text{ on } \partial M.$$

Sketlov eigenvalue

Consider the first eigenvalue of $B_g : C^\infty(\partial M) \rightarrow C^\infty(\partial M)$, where

$$B_g(f) = DN(f) + aH_g f$$

where a is constant.

Sketlov eigenvalue

Consider the first eigenvalue of $B_g : C^\infty(\partial M) \rightarrow C^\infty(\partial M)$, where

$$B_g(f) = DN(f) + aH_g f$$

where a is constant.

Theorem (H.-Shin)

Along the unnormalized CMCF

$$\frac{\partial}{\partial t} g(t) = -H_{g(t)} g(t) \text{ on } \partial M \quad \text{and} \quad R_{g(t)} = 0 \text{ in } M,$$

Sketlov eigenvalue

Consider the first eigenvalue of $B_g : C^\infty(\partial M) \rightarrow C^\infty(\partial M)$, where

$$B_g(f) = DN(f) + aH_g f$$

where a is constant.

Theorem (H.-Shin)

Along the unnormalized CMCF

$$\frac{\partial}{\partial t} g(t) = -H_{g(t)} g(t) \text{ on } \partial M \quad \text{and} \quad R_{g(t)} = 0 \text{ in } M,$$

the first eigenvalue of $B_{g(t)}$ is non-decreasing.

Steklov eigenvalue

Consider the first eigenvalue of $B_g : C^\infty(\partial M) \rightarrow C^\infty(\partial M)$, where

$$B_g(f) = DN(f) + aH_g f$$

where a is constant.

Theorem (H.-Shin)

Along the unnormalized CMCF

$$\frac{\partial}{\partial t} g(t) = -H_{g(t)} g(t) \text{ on } \partial M \quad \text{and} \quad R_{g(t)} = 0 \text{ in } M,$$

the first eigenvalue of $B_{g(t)}$ is non-decreasing.

Using CMCF, we also obtained some estimates of the first nonzero Steklov eigenvalue of (M, g) .

Yamabe soliton—the flow point of view

Suppose that M is without boundary, possibly noncompact.

Yamabe soliton—the flow point of view

Suppose that M is without boundary, possibly noncompact.
The Yamabe soliton is self-similar solution to Yamabe flow.

Yamabe soliton—the flow point of view

Suppose that M is without boundary, possibly noncompact.

The Yamabe soliton is self-similar solution to Yamabe flow.

More precisely, the **Yamabe soliton** is the solution $g(t)$ of the (unnormalized) Yamabe flow:

$$\frac{\partial}{\partial t}g(t) = -R_{g(t)}g(t)$$

and

Yamabe soliton—the flow point of view

Suppose that M is without boundary, possibly noncompact.

The Yamabe soliton is self-similar solution to Yamabe flow.

More precisely, the **Yamabe soliton** is the solution $g(t)$ of the (unnormalized) Yamabe flow:

$$\frac{\partial}{\partial t}g(t) = -R_{g(t)}g(t)$$

and

$$g(t) = \sigma(t)\psi_t^*(g_0)$$

where $\sigma : [0, \infty) \rightarrow (0, \infty)$ with $\sigma(0) = 1$,

and $\psi_t : M \rightarrow M$ is a family of diffeomorphisms with $\psi_0 = id_M$.

Yamabe soliton—the flow point of view

Suppose that M is without boundary, possibly noncompact.

The Yamabe soliton is self-similar solution to Yamabe flow.

More precisely, the **Yamabe soliton** is the solution $g(t)$ of the (unnormalized) Yamabe flow:

$$\frac{\partial}{\partial t}g(t) = -R_{g(t)}g(t)$$

and

$$g(t) = \sigma(t)\psi_t^*(g_0)$$

where $\sigma : [0, \infty) \rightarrow (0, \infty)$ with $\sigma(0) = 1$,

and $\psi_t : M \rightarrow M$ is a family of diffeomorphisms with $\psi_0 = id_M$.

This is **Yamabe soliton from flow point of view**.

Yamabe soliton—the equation point of view

We can define the Yamabe soliton from equation point of view.

Yamabe soliton—the equation point of view

We can define the Yamabe soliton from equation point of view.

Differentiating $g(t) = \sigma(t)\psi_t^*(g_0)$ with respect to t gives

$$-R_{g(t)}g(t) = \frac{\partial}{\partial t}g(t) = \dot{\sigma}(t)\psi_t^*(g_0) + \sigma(t)\frac{\partial}{\partial t}\psi_t^*(g_0).$$

Yamabe soliton—the equation point of view

We can define the Yamabe soliton from equation point of view.

Differentiating $g(t) = \sigma(t)\psi_t^*(g_0)$ with respect to t gives

$$-R_{g(t)}g(t) = \frac{\partial}{\partial t}g(t) = \dot{\sigma}(t)\psi_t^*(g_0) + \sigma(t)\frac{\partial}{\partial t}\psi_t^*(g_0).$$

Evaluating it at $t = 0$, using $\psi_0 = id_M$ and $\sigma(0) = 1$, we get

$$-R_{g_0}g_0 = \dot{\sigma}(0)g_0 + L_Xg_0,$$

where L_X is the Lie derivative with respect to X , and X is the vector field generated by ψ_t .

Yamabe soliton—the equation point of view

If we write $\lambda = -\dot{\sigma}(0)$, then

$$(\lambda - R_{g_0})g_0 = L_X g_0.$$

This is the **Yamabe soliton from the equation point of view**.

Yamabe soliton—the equation point of view

If we write $\lambda = -\dot{\sigma}(0)$, then

$$(\lambda - R_{g_0})g_0 = L_X g_0.$$

This is the **Yamabe soliton from the equation point of view**.

If the vector field is gradient, i.e. $X = \frac{1}{2}\nabla_{g_0} f$ for some function f , then

$$(\lambda - R_{g_0})g_0 = \nabla_{g_0}^2 f,$$

where $\nabla_{g_0}^2 f$ is the Hessian of f . This is called **gradient Yamabe soliton**.

Yamabe soliton—the equation point of view

If we write $\lambda = -\dot{\sigma}(0)$, then

$$(\lambda - R_{g_0})g_0 = L_X g_0.$$

This is the **Yamabe soliton from the equation point of view**.

If the vector field is gradient, i.e. $X = \frac{1}{2}\nabla_{g_0} f$ for some function f , then

$$(\lambda - R_{g_0})g_0 = \nabla_{g_0}^2 f,$$

where $\nabla_{g_0}^2 f$ is the Hessian of f . This is called **gradient Yamabe soliton**.

Fact: The Yamabe soliton from the flow point of view is equivalent to the Yamabe soliton from the equation point of view.

Theorem (di Cerbo-Disconzi)

Any compact Yamabe soliton (from flow point of view) must have constant scalar curvature.

Theorem (di Cerbo-Disconzi)

Any compact Yamabe soliton (from flow point of view) must have constant scalar curvature.

Theorem (Daskalopoulos-Sesum, Hsu)

Any compact, gradient Yamabe soliton (from equation point of view) must have constant scalar curvature.

Theorem (di Cerbo-Disconzi)

Any compact Yamabe soliton (from flow point of view) must have constant scalar curvature.

Theorem (Daskalopoulos-Sesum, Hsu)

Any compact, gradient Yamabe soliton (from equation point of view) must have constant scalar curvature.

Theorem (Daskalopoulos-Sesum)

Suppose that (M, g_0) is a gradient Yamabe soliton (from equation point of view) which is locally conformally flat and has positive sectional curvature. Then (M, g_0) has a warped-product structure.

Suppose that M has boundary ∂M , possibly noncompact.

Yamabe soliton with boundary—the flow point of view

Suppose that M has boundary ∂M , possibly noncompact.

The **Yamabe soliton with boundary** is the solution $g(t)$ of the (unnormalized) Yamabe flow with boundary:

$$\frac{\partial}{\partial t}g(t) = -R_{g(t)}g(t) \text{ in } M \quad \text{and} \quad H_{g(t)} = 0 \text{ on } \partial M$$

Yamabe soliton with boundary—the flow point of view

Suppose that M has boundary ∂M , possibly noncompact.

The **Yamabe soliton with boundary** is the solution $g(t)$ of the (unnormalized) Yamabe flow with boundary:

$$\frac{\partial}{\partial t}g(t) = -R_{g(t)}g(t) \text{ in } M \quad \text{and} \quad H_{g(t)} = 0 \text{ on } \partial M$$

and

$$g(t) = \sigma(t)\psi_t^*(g_0)$$

where $\sigma : [0, \infty) \rightarrow (0, \infty)$ with $\sigma(0) = 1$,

and $\psi_t : M \rightarrow M$ is a family of diffeomorphisms with $\psi_0 = id_M$.

Conformal mean curvature soliton—the flow point of view

The **conformal mean curvature soliton** is the solution $g(t)$ of the (unnormalized) conformal mean curvature flow:

$$\frac{\partial}{\partial t}g(t) = -H_{g(t)}g(t) \text{ on } \partial M \quad \text{and} \quad R_{g(t)} = 0 \text{ in } M$$

Conformal mean curvature soliton—the flow point of view

The **conformal mean curvature soliton** is the solution $g(t)$ of the (unnormalized) conformal mean curvature flow:

$$\frac{\partial}{\partial t}g(t) = -H_{g(t)}g(t) \text{ on } \partial M \quad \text{and} \quad R_{g(t)} = 0 \text{ in } M$$

and

$$g(t) = \sigma(t)\psi_t^*(g_0)$$

where $\sigma : [0, \infty) \rightarrow (0, \infty)$ with $\sigma(0) = 1$,

and $\psi_t : M \rightarrow M$ is a family of diffeomorphisms with $\psi_0 = id_M$.

Conformal mean curvature soliton—the flow point of view

The **conformal mean curvature soliton** is the solution $g(t)$ of the (unnormalized) conformal mean curvature flow:

$$\frac{\partial}{\partial t}g(t) = -H_{g(t)}g(t) \text{ on } \partial M \quad \text{and} \quad R_{g(t)} = 0 \text{ in } M$$

and

$$g(t) = \sigma(t)\psi_t^*(g_0)$$

where $\sigma : [0, \infty) \rightarrow (0, \infty)$ with $\sigma(0) = 1$,

and $\psi_t : M \rightarrow M$ is a family of diffeomorphisms with $\psi_0 = id_M$.

These are the Yamabe soliton with boundary and the conformal mean curvature soliton from flow point of view.

From the flow point of view, we have

Theorem (Chen-H.)

Any compact Yamabe soliton with boundary must have constant scalar curvature in M (and vanishing mean curvature on ∂M).

From the flow point of view, we have

Theorem (Chen-H.)

Any compact Yamabe soliton with boundary must have constant scalar curvature in M (and vanishing mean curvature on ∂M).

Theorem (Chen-H.)

Any compact conformal mean curvature soliton must have constant mean curvature on ∂M (and vanishing scalar curvature in M).

Proof: Since the (unnormalized) Yamabe flow with boundary preserves conformal structure, $g(t) = u(t)^{\frac{4}{n-2}} g_0$.

Proof: Since the (unnormalized) Yamabe flow with boundary preserves conformal structure, $g(t) = u(t)^{\frac{4}{n-2}} g_0$. Then

$$\frac{\partial}{\partial t} u(t) = -\frac{n-2}{4} R_{g(t)} u(t) \text{ in } M \quad \text{and} \quad \frac{\partial u(t)}{\partial \nu_{g_0}} = 0 \text{ on } \partial M.$$

Proof: Since the (unnormalized) Yamabe flow with boundary preserves conformal structure, $g(t) = u(t)^{\frac{4}{n-2}} g_0$. Then

$$\frac{\partial}{\partial t} u(t) = -\frac{n-2}{4} R_{g(t)} u(t) \text{ in } M \quad \text{and} \quad \frac{\partial u(t)}{\partial \nu_{g_0}} = 0 \text{ on } \partial M.$$

Therefore,

$$\begin{aligned} \frac{d}{dt} \left(\frac{\int_M R_{g(t)} dV_{g(t)}}{\left(\int_M dV_{g(t)}\right)^{\frac{n-2}{n}}} \right) &= \frac{d}{dt} E_{g_0}(u(t)) \\ &= -\frac{n-2}{2} \cdot \frac{\left(\int_M R_{g(t)}^2 dV_{g(t)}\right) \left(\int_M dV_{g(t)}\right) - \left(\int_M R_{g(t)} dV_{g(t)}\right)^2}{\left(\int_M dV_{g(t)}\right)^{\frac{n-2}{n}+1}} \end{aligned}$$

Cauchy-Schwarz inequality $\Rightarrow \frac{d}{dt} \left(\frac{\int_M R_{g(t)} dV_{g(t)}}{\left(\int_M dV_{g(t)}\right)^{\frac{n-2}{n}}} \right) \leq 0$, and equality holds if and only if $R_{g(t)}$ is constant.

Cauchy-Schwarz inequality $\Rightarrow \frac{d}{dt} \left(\frac{\int_M R_{g(t)} dV_{g(t)}}{\left(\int_M dV_{g(t)}\right)^{\frac{n-2}{n}}} \right) \leq 0$, and equality holds if and only if $R_{g(t)}$ is constant.

On the other hand, if $g(t) = \sigma(t)\psi_t^*(g_0)$, we have

$$\frac{\int_M R_{g(t)} dV_{g(t)}}{\left(\int_M dV_{g(t)}\right)^{\frac{n-2}{n}}} = \frac{\int_M R_{\psi_t^*(g_0)} dV_{\psi_t^*(g_0)}}{\left(\int_M dV_{\psi_t^*(g_0)}\right)^{\frac{n-2}{n}}} = \frac{\int_M R_{g_0} dV_{g_0}}{\left(\int_M dV_{g_0}\right)^{\frac{n-2}{n}}}. \quad \blacksquare$$

The equation point of view

Goal: Define the Yamabe soliton with boundary and the conformal mean curvature soliton from equation point of view.

The equation point of view

Goal: Define the Yamabe soliton with boundary and the conformal mean curvature soliton from equation point of view.

The **Yamabe soliton with boundary from equation point of view** is defined as

$$(\lambda - R_{g_0})g_0 = L_X g_0 \text{ in } M \quad \text{and} \quad H_{g_0} = 0, \langle X, \nu_{g_0} \rangle = 0 \text{ on } \partial M.$$

The equation point of view

Goal: Define the Yamabe soliton with boundary and the conformal mean curvature soliton from equation point of view.

The **Yamabe soliton with boundary from equation point of view** is defined as

$$(\lambda - R_{g_0})g_0 = L_X g_0 \text{ in } M \quad \text{and} \quad H_{g_0} = 0, \langle X, \nu_{g_0} \rangle = 0 \text{ on } \partial M.$$

The **conformal mean curvature soliton from equation point of view** is defined as

$$R_{g_0} = 0 \text{ in } M \quad \text{and} \quad (\lambda - H_{g_0})g_0 = L_X g_0, \langle X, \nu_{g_0} \rangle = 0 \text{ on } \partial M.$$

The equation point of view

Theorem (H.-Shin)

The Yamabe soliton with boundary from flow point of view is equivalent to the Yamabe soliton with boundary from equation point of view.

The equation point of view

Theorem (H.-Shin)

The Yamabe soliton with boundary from flow point of view is equivalent to the Yamabe soliton with boundary from equation point of view.

Theorem (H.-Shin)

The conformal mean curvature soliton from flow point of view is equivalent to the conformal mean curvature soliton from equation point of view.

The equation point of view

Example: Consider (\mathbb{R}_+^n, g_e) .

The equation point of view

Example: Consider (\mathbb{R}_+, g_e) . Then $R_{g_e} = 0$ in \mathbb{R}_+^n and $H_{g_e} = 0$ on $\partial\mathbb{R}_+^n$.

The equation point of view

Example: Consider (\mathbb{R}_+^n, g_e) . Then $R_{g_e} = 0$ in \mathbb{R}_+^n and $H_{g_e} = 0$ on $\partial\mathbb{R}_+^n$. Also, $\nu_{g_e} = (0, 0, \dots, 0, -1)$.

The equation point of view

Example: Consider (\mathbb{R}_+^n, g_e) . Then $R_{g_e} = 0$ in \mathbb{R}_+^n and $H_{g_e} = 0$ on $\partial\mathbb{R}_+^n$. Also, $\nu_{g_e} = (0, 0, \dots, 0, -1)$.

Thus, if we choose $f = f(x_1, \dots, x_n) = a_1x_1 + \dots + a_{n-1}x_{n-1}$, then $X = \frac{1}{2}\nabla_{g_e} f = (\frac{a_1}{2}, \dots, \frac{a_{n-1}}{2}, 0)$

The equation point of view

Example: Consider (\mathbb{R}_+^n, g_e) . Then $R_{g_e} = 0$ in \mathbb{R}_+^n and $H_{g_e} = 0$ on $\partial\mathbb{R}_+^n$. Also, $\nu_{g_e} = (0, 0, \dots, 0, -1)$.

Thus, if we choose $f = f(x_1, \dots, x_n) = a_1x_1 + \dots + a_{n-1}x_{n-1}$, then $X = \frac{1}{2}\nabla_{g_e} f = (\frac{a_1}{2}, \dots, \frac{a_{n-1}}{2}, 0)$

$$\Rightarrow L_X g_e = \nabla_{g_e}^2 f = 0 \quad \text{and} \quad \langle X, \nu_{g_e} \rangle = 0.$$

The equation point of view

Example: Consider (\mathbb{R}_+^n, g_e) . Then $R_{g_e} = 0$ in \mathbb{R}_+^n and $H_{g_e} = 0$ on $\partial\mathbb{R}_+^n$. Also, $\nu_{g_e} = (0, 0, \dots, 0, -1)$.

Thus, if we choose $f = f(x_1, \dots, x_n) = a_1 x_1 + \dots + a_{n-1} x_{n-1}$, then $X = \frac{1}{2} \nabla_{g_e} f = (\frac{a_1}{2}, \dots, \frac{a_{n-1}}{2}, 0)$

$$\Rightarrow L_X g_e = \nabla_{g_e}^2 f = 0 \quad \text{and} \quad \langle X, \nu_{g_e} \rangle = 0.$$

Therefore, we have

$$-R_{g_e} g_e = L_X g_e \text{ in } \mathbb{R}_+^n \quad \text{and} \quad H_{g_e} = 0, \langle X, \nu_{g_e} \rangle = 0 \text{ on } \partial\mathbb{R}_+^n$$

i.e. (\mathbb{R}_+^n, g_e) is a (steady, nontrivial) gradient Yamabe soliton with boundary (from equation point of view).

The equation point of view

Example: Consider (\mathbb{R}_+^n, g_e) . Then $R_{g_e} = 0$ in \mathbb{R}_+^n and $H_{g_e} = 0$ on $\partial\mathbb{R}_+^n$. Also, $\nu_{g_e} = (0, 0, \dots, 0, -1)$.

Thus, if we choose $f = f(x_1, \dots, x_n) = a_1 x_1 + \dots + a_{n-1} x_{n-1}$, then $X = \frac{1}{2} \nabla_{g_e} f = (\frac{a_1}{2}, \dots, \frac{a_{n-1}}{2}, 0)$

$$\Rightarrow L_X g_e = \nabla_{g_e}^2 f = 0 \quad \text{and} \quad \langle X, \nu_{g_e} \rangle = 0.$$

Therefore, we have

$$-R_{g_e} g_e = L_X g_e \text{ in } \mathbb{R}_+^n \quad \text{and} \quad H_{g_e} = 0, \langle X, \nu_{g_e} \rangle = 0 \text{ on } \partial\mathbb{R}_+^n$$

i.e. (\mathbb{R}_+^n, g_e) is a (steady, nontrivial) gradient Yamabe soliton with boundary (from equation point of view).

Similarly, (\mathbb{R}_+^n, g_e) is a (steady, nontrivial) gradient conformal mean curvature soliton (from equation point of view).

Theorem (H.-Shin)

Any compact gradient Yamabe soliton with boundary from equation point of view must have constant scalar curvature in M (and vanishing mean curvature on ∂M).

The equation point of view

Theorem (H.-Shin)

Any compact gradient Yamabe soliton with boundary from equation point of view must have constant scalar curvature in M (and vanishing mean curvature on ∂M).

Theorem (H.-Shin)

Any compact gradient conformal mean curvature soliton from equation point of view must have constant mean curvature on ∂M (and vanishing scalar curvature in M).

The equation point of view

Proof: If (M, g_0, f) is a gradient Yamabe soliton with boundary, then

$$\nabla_{g_0}^2 f = (\lambda - R_{g_0})g_0 \text{ in } M \quad \text{and} \quad \frac{\partial f}{\partial \nu_{g_0}} = 0 \text{ on } \partial M.$$

The equation point of view

Proof: If (M, g_0, f) is a gradient Yamabe soliton with boundary, then

$$\nabla_{g_0}^2 f = (\lambda - R_{g_0})g_0 \text{ in } M \quad \text{and} \quad \frac{\partial f}{\partial \nu_{g_0}} = 0 \text{ on } \partial M.$$

Taking trace $\Rightarrow \Delta_{g_0} f = n(\lambda - R_{g_0})$ in M .

The equation point of view

Proof: If (M, g_0, f) is a gradient Yamabe soliton with boundary, then

$$\nabla_{g_0}^2 f = (\lambda - R_{g_0})g_0 \text{ in } M \quad \text{and} \quad \frac{\partial f}{\partial \nu_{g_0}} = 0 \text{ on } \partial M.$$

Taking trace $\Rightarrow \Delta_{g_0} f = n(\lambda - R_{g_0})$ in M . Also, one has

$$(n-1)\Delta_{g_0} R_{g_0} = \frac{1}{2}\langle \nabla_{g_0} R_{g_0}, \nabla_{g_0} f \rangle + R_{g_0}(\lambda - R_{g_0}) \text{ in } M, \quad \frac{\partial R_{g_0}}{\partial \nu_{g_0}} = 0 \text{ on } \partial M.$$

The equation point of view

Proof: If (M, g_0, f) is a gradient Yamabe soliton with boundary, then

$$\nabla_{g_0}^2 f = (\lambda - R_{g_0})g_0 \text{ in } M \quad \text{and} \quad \frac{\partial f}{\partial \nu_{g_0}} = 0 \text{ on } \partial M.$$

Taking trace $\Rightarrow \Delta_{g_0} f = n(\lambda - R_{g_0})$ in M . Also, one has

$$(n-1)\Delta_{g_0} R_{g_0} = \frac{1}{2}\langle \nabla_{g_0} R_{g_0}, \nabla_{g_0} f \rangle + R_{g_0}(\lambda - R_{g_0}) \text{ in } M, \quad \frac{\partial R_{g_0}}{\partial \nu_{g_0}} = 0 \text{ on } \partial M.$$

Integrating the first equation over M , we get

$$(n-1) \int_M \Delta_{g_0} R_{g_0} = \frac{1}{2} \int_M \langle \nabla_{g_0} R_{g_0}, \nabla_{g_0} f \rangle + \int_M R_{g_0}(\lambda - R_{g_0}),$$

The equation point of view

Proof: We get

$$0 = \left(1 - \frac{n}{2}\right) \int_M R_{g_0} (\lambda - R_{g_0}). \quad (1)$$

The equation point of view

Proof: We get

$$0 = \left(1 - \frac{n}{2}\right) \int_M R_{g_0} (\lambda - R_{g_0}). \quad (1)$$

Finally, integrating $\Delta_{g_0} f = n(\lambda - R_{g_0})$ over M and using $\frac{\partial f}{\partial \nu_{g_0}} = 0$ on ∂M , we obtain

$$0 = \int_M (\lambda - R_{g_0}) \quad (2)$$

The equation point of view

Proof: We get

$$0 = \left(1 - \frac{n}{2}\right) \int_M R_{g_0} (\lambda - R_{g_0}). \quad (1)$$

Finally, integrating $\Delta_{g_0} f = n(\lambda - R_{g_0})$ over M and using $\frac{\partial f}{\partial \nu_{g_0}} = 0$ on ∂M , we obtain

$$0 = \int_M (\lambda - R_{g_0}) \quad (2)$$

Combining (1) and (2), we get

$$\int_M (\lambda - R_{g_0})^2 = \int_M R_{g_0} (\lambda - R_{g_0}) - \lambda \int_M (\lambda - R_{g_0}) = 0,$$

which gives $R_{g_0} \equiv \lambda$.

The equation point of view

Theorem (H.-Shin)

Suppose that (M, g_0) is a Yamabe soliton with boundary from equation point of view. If (M, g_0) is locally conformal flat and has positive sectional curvature,

The equation point of view

Theorem (H.-Shin)

Suppose that (M, g_0) is a Yamabe soliton with boundary from equation point of view. If (M, g_0) is locally conformal flat and has positive sectional curvature, then

$$(M, g_0) = (\mathbb{R}, dr^2) \times_{w(r)} (\mathbb{S}_+^{n-1}, g_{\mathbb{S}_+^{n-1}}).$$

$$\dim M = 2$$

Suppose M is a 2-dimensional manifold with boundary ∂M .

Suppose M is a 2-dimensional manifold with boundary ∂M .

The (unnormalized) **Gauss curvature flow with boundary** is

$$\frac{\partial}{\partial t} g(t) = -K_{g(t)} g(t) \text{ in } M \quad \text{and} \quad k_{g(t)} = 0 \text{ on } M.$$

Here, $K_{g(t)}$ is the Gauss curvature, and $k_{g(t)}$ is the geodesic curvature.

Suppose M is a 2-dimensional manifold with boundary ∂M .

The (unnormalized) **Gauss curvature flow with boundary** is

$$\frac{\partial}{\partial t} g(t) = -K_{g(t)} g(t) \text{ in } M \quad \text{and} \quad k_{g(t)} = 0 \text{ on } M.$$

Here, $K_{g(t)}$ is the Gauss curvature, and $k_{g(t)}$ is the geodesic curvature.

The Gauss curvature flow with boundary is a two-dimensional version of the Yamabe flow with boundary.

Suppose M is a 2-dimensional manifold with boundary ∂M .

The (unnormalized) **geodesic curvature flow** is

$$\frac{\partial}{\partial t}g(t) = -k_{g(t)}g(t) \text{ on } \partial M \quad \text{and} \quad K_{g(t)} = 0 \text{ in } M.$$

Suppose M is a 2-dimensional manifold with boundary ∂M .

The (unnormalized) **geodesic curvature flow** is

$$\frac{\partial}{\partial t} g(t) = -k_{g(t)} g(t) \text{ on } \partial M \quad \text{and} \quad K_{g(t)} = 0 \text{ in } M.$$

The geodesic curvature flow is a two-dimensional version of the conformal mean curvature flow.

$$\dim M = 2$$

We have defined the **Gauss curvature soliton with boundary**, which is the self-similar solution of the Gauss curvature flow with boundary.

$$\dim M = 2$$

We have defined the **Gauss curvature soliton with boundary**, which is the self-similar solution of the Gauss curvature flow with boundary.

Also, we have defined the **geodesic curvatures soliton**, which is the self-similar solution of the geodesic curvature flow.

We have defined the **Gauss curvature soliton with boundary**, which is the self-similar solution of the Gauss curvature flow with boundary.

Also, we have defined the **geodesic curvatures soliton**, which is the self-similar solution of the geodesic curvature flow.

Theorem (H.-Shin)

(i) *Any compact Gauss curvature soliton with boundary must have constant Gauss curvature in M .*

We have defined the **Gauss curvature soliton with boundary**, which is the self-similar solution of the Gauss curvature flow with boundary.

Also, we have defined the **geodesic curvatures soliton**, which is the self-similar solution of the geodesic curvature flow.

Theorem (H.-Shin)

- (i) Any compact Gauss curvature soliton with boundary must have constant Gauss curvature in M .
- (ii) Any compact geodesic curvature soliton must have constant geodesic curvature in M .

- ▶ (Joint with Jinwoo Shin) Slowly converging Yamabe-type flow on manifolds with boundary. *Commun. Contemp. Math.* (2022), accepted.
- ▶ (Joint with Jinwoo Shin) Yamabe solitons with boundary. *Ann. Mat. Pura Appl. (4)* **202** (2023), no. 5, 2219–2253.
- ▶ (Joint with Jinwoo Shin) Evolution of the Steklov eigenvalue along the conformal mean curvature flow. *J. Geom. Phys.* **173** (2022), Paper No. 104436.
- ▶ (Joint with Xuezhong Chen) Conformal curvature flows on compact manifold of negative Yamabe constant. *Indiana U. Math. J.* **67** (2018), 537–581.

Thank you very much for your attention!