On the Rokhlin type formulae of Spin geometry

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- Introduce the Spin characteristic classes; and
- Present an application of these classes to Spin geometry.

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Arrangement of the talk:

- The Spin characteristic classes;
- The classical Rokhlin formula;
- The Rokhlin type formulae (in the Spin characteristic classes);
- Applications to Spin geometry.

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- Haefliger (1956) found that the second Stiefel Whitney class $w_2(M)$ is the only obstruction to the existence of a spin structure on an orientable Riemannian manifold M.
- This was extended by Borel and Hirzebruch (1958) to cases of vector bundles, and by Karoubi (1968) to the non-orientable pseudo-Riemannian manifolds.

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Assume that the ring $H^*(B_{Spin(n)})$ has been determined. We can specify a minimal set $\{q_1, \dots, q_r\}$ of generators of the ring $H^*(B_{Spin(n)})$. We can then introduce **the spin characteristic classes** for a spin vector bundle ξ over a space X with classifying map $f : X \to B_{Spin(n)}$ by setting

$$q_i(\xi) = f^*(q_i) \in H^*(X), \ 1 \le i \le r.$$

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Remark. The classical Chern characteristic classes for the complex vector bundles, and the Pontryagin characteristic classes for the oriented real vector bundles can be defined in the same way.

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- The mod 2 cohomology of the space $B_{Spin(n)}$ was computed by Borel (1953) for $n \le 10$, and was completed by Quillen (1972) for all n.
- Thomas (1962) calculated the integral cohomology of B_{Spin(∞)} in the stable range, but the result was subject to the choice of two sets {Φ_i, Ψ_i} of indeterminants.

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- In the context of Weyl invariants, a description of the integral cohomology H*(B_{Spin(n)}) was formulated by Benson and Wood (1995), where explicit generators and relations are absent:

"We have not set about the rather daunting task of using this description to give explicit generators and relations"

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for the respective algebras $H^*(B;Z)$ and $H^*(B;Z_2)$. Set

$$D^{i} = D \cap H^{i}(B;Z), D_{2}^{i} = D_{2} \cap H^{i}(B;Z_{2}) \quad (i \ge 0).$$

We shall prove

THEOREM (1.2). There are cohomology classes $\{Q_i\}, \{\Phi_i\}, \{\Psi_i\}$ $(i \ge 1)$ with the **以**後日 日前日 following properties:

(1.3) $Q_i \in H^{4i}(\hat{B},Z), \Phi_i \in D^{4i} \subset H^{4i}(B;Z), \quad \Psi_i \in D_2^i \subset H^i(B;Z_2).$

(1.4) If i is not a power of 2, then

 $Q_i = \pi^* P_i, \quad \Phi_i = 0, \quad \Psi_i = 0.$

Let $j = 2^r$, for $r = 0, 1, \cdots$. Then

- $\pi^* P_{2j} = 2Q_{2j} + Q_j^2 \pi^* \Phi_{2j}, \quad \pi^* P_1 = 2Q_1;$ (1.5) $\rho_2(Q_i) = \pi^*(W_{4i} + \Psi_{4i}), \qquad \rho_2(\Phi_i) = (\Psi_{2i})^2.$
- (1.6)

Moreover.

(1.7)
$$H^*(\hat{B};Z) = Z[Q_1, Q_2, \cdots] \oplus \hat{T}, \text{ where } 2\hat{T} = 0.$$

Furthermore, if $\{Q'_i\}(i > 1)$ is a second set of cohomology classes satisfying

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11. PIECING TOGETHER THE INTEGRAL COHOMOLOGY

We calculated $H^*(BSpin(n); \mathbb{Z})/torsion$ in the last section. We saw that the torsion is killed by multiplication by two, and so we have a pullback diagram

> $H^*(B\operatorname{Spin}(n); \mathbb{Z}) \rightarrow H^*(B\operatorname{Spin}(n); \mathbb{Z})/\operatorname{torsion}$ $\downarrow \qquad \qquad \downarrow \qquad \qquad \downarrow$ Kerß

Thus we have the following.

THEOREM 11.1. $H^*(B \operatorname{Spin}(n); \mathbb{Z})$ is isomorphic to the subring of

Ker $\beta \oplus H^*(BSpin(n); \mathbb{Z})/torsion$

consisting of pairs of elements with the same image in H_e.

We have not set about the rather daunting task of using this description to give explicit generators and relations.

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Theorem (Duan, 2018). The integral cohomology of B_{Spin} has the presentation

$$H^*(B_{Spin}) = \mathbb{Z}[Q_1, Q_2, Q_3, \cdots] \oplus \pi^* \tau(B_{SO})$$

where the generators Q_k , $degQ_k = 4k$, are *characterized uniquely* by the following properties:

i) if
$$k > 1$$
 is not a power of 2, then $Q_k = \pi^* p_k$.
ii) if $k = 2^r$ then
 $\rho_2(Q_k) = \pi^*(w_2^{(k+1)}), 2Q_{2k} + Q_k^2 = \pi^* f(w_2^{(k+1)}).$

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Definition. For a Spin bundle ξ over a space X with classifying map $f: X \to B_{Spin}$, the k^{th} Spin classes of ξ is

$$q_k(\xi):=f^*(Q_k),\ k\geq 1.$$

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 $H^{2m}(M) \times H^{2m}(M) \rightarrow H^{4m}(M) = \mathbb{Z}$

is an integral unimodular symmetric matrix A, whose signature will be denoted by σ_M .

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Theorem (Rokhlin, 1952). The signature σ_M of a 4-dimensional smooth spin manifold M satisfies that $\sigma_M \equiv 0 \mod 2^4$.

Remark. These results motivated the classifications on simply connected 4 dimensional manifolds due to Freedman **(Topological)** and Donaldson **(smooth)** in the 1980's.

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In addition, the relations on the ring $H^*(Spin(n))$ suffices for us to express the Pontryagin classes by the Spin classes, such as

$$p_1 = 2q_1, p_2 = 2q_2 - q_1^2, p_3 = q_3, p_4 = 2q_4 - q_2^2, \cdots$$

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Let M^{4m} be a 4m dimensional smooth manifold with total Pontryagin class $1 + p_1 + \cdots + p_m$. Recall that the \widehat{A} genus α_m and the L genus τ_m of M^{4m} can be expressed as polynomials in the p_1, \cdots, p_m of the forms

$$lpha_m = a_m \cdot p_m + I_m(p_1, \cdots, p_{m-1})$$
 and
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Haibao Duan (CAS)

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where a_m and b_m are certain non-zero rationals. These allow us to eliminate p_m to obtain the following expression of τ_m without involving p_m :

Rokhlin type formula in Pontryagin classes

$$\tau_m = \frac{b_m}{a_m}(\alpha_m - I_m(p_1, \cdots, p_{m-1})) + k_m(p_1, \cdots, p_{m-1}).$$

For $1 \leq m \leq$ 4 the polynomials α_m and τ_m are

$$\begin{aligned} \alpha_{1} &= -\frac{1}{24} p_{1}; \\ \alpha_{2} &= \frac{1}{2^{7} \cdot 3^{2} \cdot 5} (-4 p_{2} + 7 p_{1}^{2}) \\ \alpha_{3} &= \frac{1}{2^{10} \cdot 3^{3} \cdot 5 \cdot 7} (-16 p_{3} + 44 p_{2} p_{1} - 31 p_{1}^{3}) \\ \alpha_{4} &= \frac{1}{2^{15} \cdot 5^{2} \cdot 3^{4} \cdot 7} (-192 p_{4} + 512 \cdot p_{1} p_{3} + 208 p_{2}^{2} - 904 p_{1}^{2} p_{2} + 381 p_{1}^{4}) \end{aligned}$$

and

$$\begin{split} \tau_1 &= \frac{1}{3} p_1; \\ \tau_2 &= \frac{1}{3^2 \cdot 5} (7 p_2 - p_1^2); \\ \tau_3 &= \frac{1}{3^3 \cdot 5 \cdot 7} (62 p_3 - 13 p_2 p_1 + 2 p_1^3); \\ \tau_4 &= \frac{1}{3^4 \cdot 5^2 \cdot 7} (381 p_4 - 71 \cdot p_1 p_3 - 19 p_2^2 + 22 p_1^2 p_2 - 3 p_1^4), \end{split}$$

respectively.

If the manifold M^{4m} is spin, we can use the Spin characteristic classes in places of p_1, \dots, p_{m-1} to get the following surprisingly simple formulae:

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These formulae are called as "**the Rokhlin type formulae**" for the following reason:

Corollary. In accordance to m = 1, 2, 3, 4, the signature σ_M of an almost parallelizable smooth manifold M^{4m} is divisible, respectively, by

$$2^4$$
, $2^5 \cdot (2^3 - 1)$, $2^8 \cdot (2^5 - 1)$, $2^9 \cdot (2^7 - 1)$.

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Remark. Atiyah-Singer Index Theorem, 1963: The \widehat{A} genus α_m of a smooth spin manifold M^{4m} is equal to the index of its Dirac operator.

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Application 1. The signature σ_M of an 8 dimensional smooth spin manifold M satisfies that

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A spin manifold is called *string* if its first Spin class q_1 vanishes. In this case the second Spin class q_2 is divisible by 3, while the formula τ_4 implies:

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Definition. For a symmetric unimodular integral matrix $A = (a_{ij})_{n \times n}$ of rank *n*, and a sequence $b = (b_1, \dots, b_n)$ of *n* integers, we say the pair (A, b) to be **a Wall pair** if the following constraints are satisfied:

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i) M - {o} admits a smooth structure for some point o ∈ M;
ii) there is a basis {x₁, ··· , x_n} on H⁴(M⁸) so that the corresponding intersection form is A;
iii) the first Spin characteristic class of M is q₁ = b₁x₁ + ··· + b_nx_n ∈ H⁴(M⁸).

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Theorem 1. Let M^8 be the manifold associated to a Wall pair (A, b). There exists a smooth structure on M^8 extending the one on $M - \{o\}$ if and only if

$$sign(A) \equiv bAb^{ au} \mod 2^5 \cdot (2^3 - 1),$$

where b^{τ} denotes the transpose of the row vector b.

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Theorem 2. If (A, b) is a Wall pair so that

 $sign(A) \neq bAb^{\tau} mod2^5 \cdot (2^3 - 1),$

then the corresponding manifold M does not admit any smooth structure. \Box

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Conversely, for those M which admit smooth structures, their tangent invariants q(M) can be determined completely.

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Conversely, for those M which admit smooth structures, their tangent invariants q(M) can be determined completely.

Theorem 3. If the manifold M associated to a Wall pair (A, b) admits a smooth structure, then its total Spin characteristic class is

$$q(M) = 1 + (b_1x_1 + \dots + b_1x_n) + \frac{45 \cdot sign(A) + 11 \cdot bAb^{\tau}}{2 \cdot (2^3 - 1)} \cdot \omega_M$$

where $2 \cdot (2^3 - 1)$ divides $45 \cdot sign(A) + 11 \cdot bAb^{\tau}$.

Theorem 4. For a manifold M associated to a Wall pair (A, b) the following statements are equivalent:

i) *M* is smooth and admits a metric with positive scalar curvature ii) $sign(A) = bAb^{\tau}$; iii) $q_2(M) = 4sign(A) \cdot \omega_M.\Box$

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Theorem 4. For a manifold M associated to a Wall pair (A, b) the following statements are equivalent:

i) M is smooth and admits a metric with positive scalar curvature
ii) sign(A) = bAb^τ;
iii) q₂(M) = 4sign(A) · ω_M.□

Therefore, the problem of finding all the 3 connected 8 dimensional smooth manifolds that have a metric with positive scalar curvature is equivalent to the number theoretical problem of finding those Wall pairs (A, b) satisfying the quadratic equation ii) above.

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Thanks!

Haibao Duan (CAS)

On the Rokhlin type formulae of Spin geome

October 16, 2018

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